

Improved Photoelectrochemical Performance of GaN Nanopillar Photoanodes

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Abstract

In this work, we report on the photoelectrochemical (PEC) investigation of n-GaN nanopillar (NP) photoanodes fabricated using metal organic chemical vapour deposition and top-down approach. Substantial improvement in photocurrents is observed for GaN NP photoanodes compared to their planar counterparts. The role of carrier concentration and NP dimensions on the PEC performance of NP photoanodes is further elucidated. Photocurrent density is almost doubled for doped NP photoanodes whereas no improvement is noticed for undoped NP photoanodes. While the diameter of GaN NP is found to influence the onset potential, carrier concentration is found to affect both the onset and overpotential of the electrodes. Optical and electrochemical impedance spectroscopy characterizations are utilised to further explain the PEC results of NP photoanodes. Finally, improvement in the photostability of NP photoanodes with the addition of NiO as a co-catalyst is investigated.

Keywords: GaN, Nanopillars, Photoanodes

Introduction:

Generation of hydrogen fuel by splitting water using renewable solar energy is a highly desirable and exciting prospect for the sustainable development of human race.[1-3] III-V nitrides have recently been gaining considerable research interest in solar water splitting owing to their outstanding properties.[4-8] The advantages of nitride alloys include tunable bandgap, high crystal quality and stability against photo and chemical corrosion. For example, the band gap of $\text{In}_x\text{Ga}_{1-x}\text{N}$ alloys can be tuned from 3.4 to 0.7 eV depending on the In composition. This facilitates the manipulation of light absorption across the entire solar region for efficient solar water splitting

while straddling the water oxidation and reduction potentials.[9] Typically, GaN layer is grown epitaxially in single crystal quality which is highly beneficial for efficient charge transfer. Moreover, the strong ionic bonding nature of III-nitride alloys helps in stabilizing the photoelectrodes against photo-corrosion. All these superior properties of III-nitrides make them excellent candidate materials for solar water splitting applications.

In addition, nanostructures offer great potential for solar water splitting over their planar counterparts. Nanostructures possess several essential attributes towards achieving efficient water splitting such as enhanced light absorption, reduced carrier transfer lengths, large surface to facilitate efficient charge transfer at semiconductor/electrolyte interface and enlarged depletion area at the nanostructured surface that drives the charge separation efficiently. [10-15] GaN nanostructures can be synthesized via bottom-up and top-down approaches. In the bottom-up approach, nanostructures are grown either by selective area epitaxial growth or vapour-liquid-solid growth [16, 17]. In the top-down approach, nanopillars (NPs) are formed by reactive ion etching of patterned GaN epilayers.[18-21] Controlling the morphology and doping of NPs is a main setback for bottom-up approaches.[22] On the other hand, NPs fabricated using the top-down approach exhibit controlled morphology with uniform doping as it allows the use of high quality GaN epilayer grown using well-matured planar growth technology wherein doping concentration is precisely controlled. The morphology of the NPs can be controlled by well-established lithography techniques or random masking techniques. Self-assembled random mask cost is effective technique for produce the wafer-scale etch mask.[23, 24] In addition, the density and size of the NP can be controlled by simply varying the thickness of Ni film. As a result, top-down approach produces nanostructures in a cost effective manner at the wafer-scale level. Despite these merits, there are very few reports on the photoelectrochemical (PEC) studies of GaN NPs fabricated using top-down approach.

Carrier concentration and dimensions of nanostructures are important parameters towards realizing efficient water splitting. There are considerable number of research reports on nitride-based PEC water splitting [25-31] and overall direct water splitting. [32-35] However, there are very few studies focused on the influence of carrier concentration and NP dimensions on solar water splitting. Therefore, it is critical to understand the dependence of PEC performance on NP dimensions and carrier concentration. J. Benton et al. [36] reported significant enhancement of

photocurrent and hydrogen generation for n-GaN/MQW(InGaN/GaN)/p-GaN NPs, fabricated by the top-down process, compared with their planar counterparts. However, the structure applied in their work includes a p-n junction and InGaN multiple quantum wells, which complicate the PEC analysis of GaN NPs. W. J. Tseng et al. [37, 38] investigated the influence of length of the NPs and surface damage to NPs that occurred during the dry etching on the PEC performance of NP photoanodes, however, the photocurrents were significantly low. Nanowires grown in the bottom-up approaches exhibit varying doping concentration along the nanowire length. This makes it difficult to understand the influence of carrier concentration on the PEC performance of NP photoanodes. Therefore, top-down approach is an ideal method for fabricating NPs to investigate the influence of carrier concentration and NP dimensions on the PEC performance of NP photoanodes.

In this work, we report on the fabrication of n-GaN NP photoanodes using the top-down approach and the investigation of the influence of carrier concentration and NP dimensions on their PEC performance. Self-assembled random-mask technique was used for producing etch mask. Impedance and optical spectroscopy measurements were carried out to explain the PEC performance of NP photoanodes. The dependence of depletion layer thickness in nanostructures on carrier concentration and NP dimensions is used to explain the shift in anodic potential of the NPs with changing dimensions and carrier concentrations.

Experimental

For the fabrication of GaN NPs, 2 μm -thick GaN epitaxial film with controlled carrier concentration was first grown on sapphire wafer using metal organic chemical vapour deposition (MOCVD). Carrier concentration of the GaN epilayers was adjusted by varying the silane flow during the growth. GaN NPs were fabricated from the epilayers using inductively coupled plasma-reactive ion etching (ICP-RIE) and self-assembled random nano-mask techniques.[24] SiO_2/Ni mask was used as an etch mask for the fabrication of GaN NPs. For the ICP-RIE etching of SiO_2 , 40 sccm of CHF_3 gas was passed through the ICP-RIE chamber that was maintained at 20 $^\circ\text{C}$ electrode temperature, 10 mT pressure and 30/1000 watt RF/ICP powers. On the other hand, GaN ICP-RIE etching has been carried out at 60 $^\circ\text{C}$ electrode temperature by passing 25/2/5 sccm of $\text{Cl}_2/\text{H}_2/\text{Ar}$ gases at a chamber pressure of 4 mT and 200/1000 watt RF/ICP powers.

NiO nanoparticle co-catalyst decoration on the surface of GaN NPs was achieved by depositing 5 nm Ni film on the GaN NPs using e-beam evaporation (E-beam, Temescal BJD-2000). This Ni film coated NPs were furnace annealed at 550 °C in air for 30 min to convert the Ni film into NiO nano-islands. Field emission scanning electron microscopy (FEI Helios Nanolab 600) was employed to investigate the morphology of GaN NPs and NiO-coated NPs. Perkin Elmer spectrophotometer equipped with an integrating sphere was used to characterize the diffuse reflectance of planar GaN and GaN NPs.

To measure the PEC properties of planar and NPs photoanodes Ni/Au (30/150nm) metal contacts were deposited onto the planar and NP GaN samples. The contacts were then subjected to rapid thermal annealing at 400 °C in argon atmosphere for 5 min. PEC measurements of these photoanodes were carried out under a solar simulator (Newport) fitted with air mass 1.5 filters in 1 M NaOH electrolyte solution in a quartz cell consisting of the GaN working electrode, a Pt counter electrode and a Ag/AgCl reference electrode.

Results and Discussion

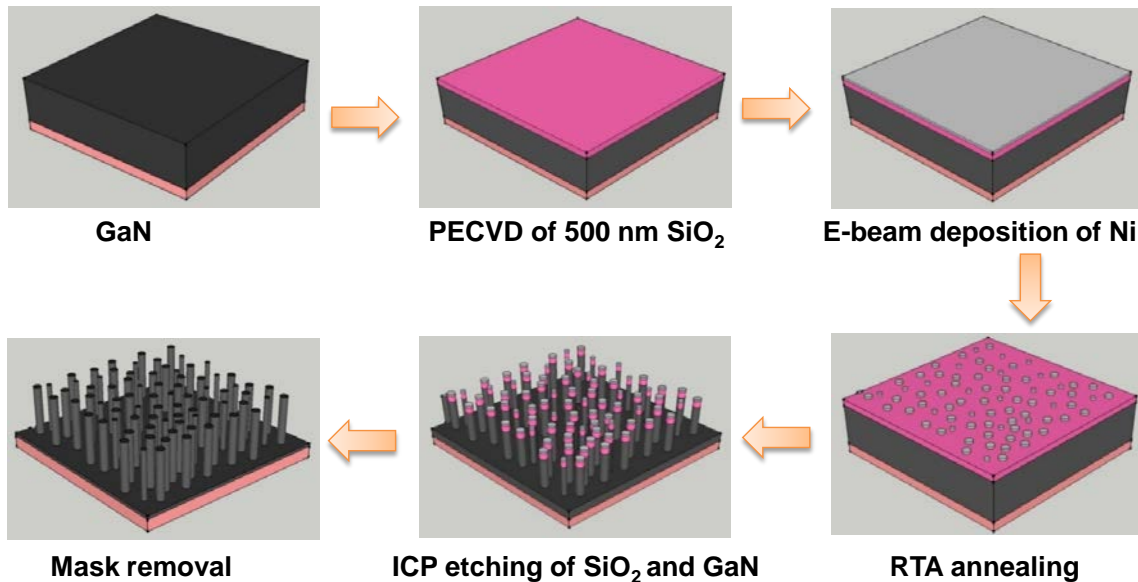


Fig 1. Illustration of the processing steps involved in the fabrication of GaN NPs using top-down approach.

Fig 1 shows the schematic illustration of top-down fabrication of GaN NPs. GaN nanopyllars were fabricated by ICP-RIE etching of SiO₂/Ni masked n-doped and undoped GaN wafers as illustrated

in Fig 1. SiO₂/Ni hard mask was obtained using plasma enhanced chemical vapour deposition (Oxford Plasmalab 100) of 500 nm-thick SiO₂ film on the GaN wafers followed by electron beam evaporation of 5-15 nm thick nickel film. The Ni film was converted into nano-islands by rapid thermal annealing at 900 °C for 2 minutes in Ar atmosphere. Next, the SiO₂/Ni random nano-pattern was created by ICP-RIE (Versaline LL) etching of SiO₂ surrounding the nickel islands. The SiO₂/Ni random mask was then used to selectively etch GaN and produce GaN NPs using ICP-RIE. After etching, GaN NPs were dipped in HF solution to remove the leftover SiO₂/Ni mask. Fig 2 shows the GaN nanopillars after removing the mask. The length and diameter of the NPs were controlled by controlling the etching time and Ni film thickness, respectively. While the etching time determines the length of the NPs, diameter of the NPs depends on the size of the Ni nano-islands. The size of the Ni nano-islands increases with increasing Ni film thickness. Fig 2 (a-c) show the SEM images of GaN NPs of lengths 200, 600 & 900 nm etched for 60, 90 and 120 s, respectively, using 5 nm Ni masked sample. Fig. 2 (d-f) show the NPs with increasing diameter fabricated using 5, 10 and 15 nm Ni mask, respectively, after etching for 120 sec. All the NPs have uniform length and smooth sidewalls irrespective of the etch time and Ni mask size. The average diameter of the NPs measured from the SEM images is 84, 144 and 202 nm for the 5, 10 and 15 nm Ni film masked samples, respectively. The diameter distribution becomes wider with increasing thickness of the Ni film (Fig SI1).

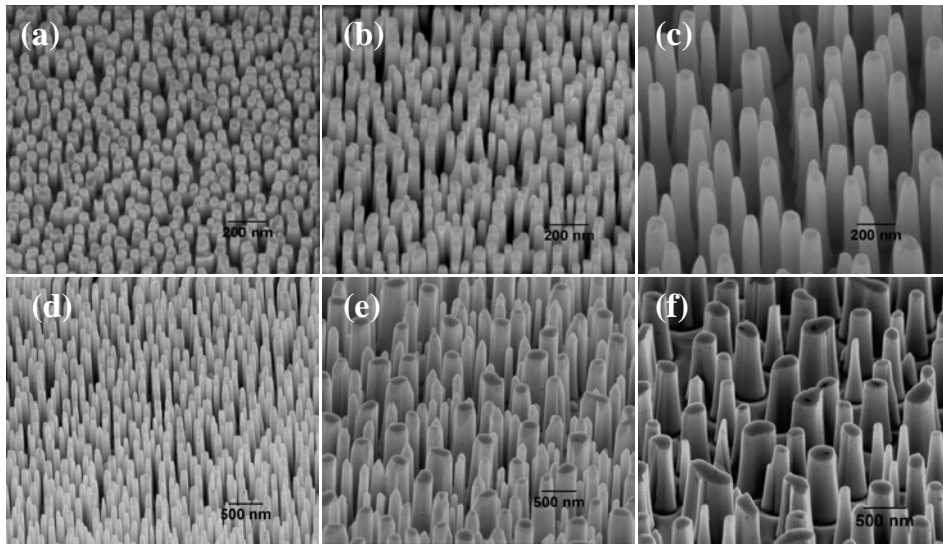


Fig 2. SEM images of the top-down fabricated GaN NPs (images took at 30° tilt angle) with increasing lengths (a-c) using 5 nm Ni mask obtained by varying the etch times from 60 to 120 s and increasing diameter (d-f) (images took at 52° tilt angle) using 5, 10 and 15 nm Ni mask etched for 120 s.

Table 1: Sample numbers and the corresponding carrier concentrations

Sample name	Carrier concentration (cm ⁻³)
Sample 1 (S1)	Undoped
Sample 2 (S2)	1 x 10 ¹⁸
Sample 3 (S3)	5 x 10 ¹⁸
Sample 4 (S4)	1 x 10 ¹⁹

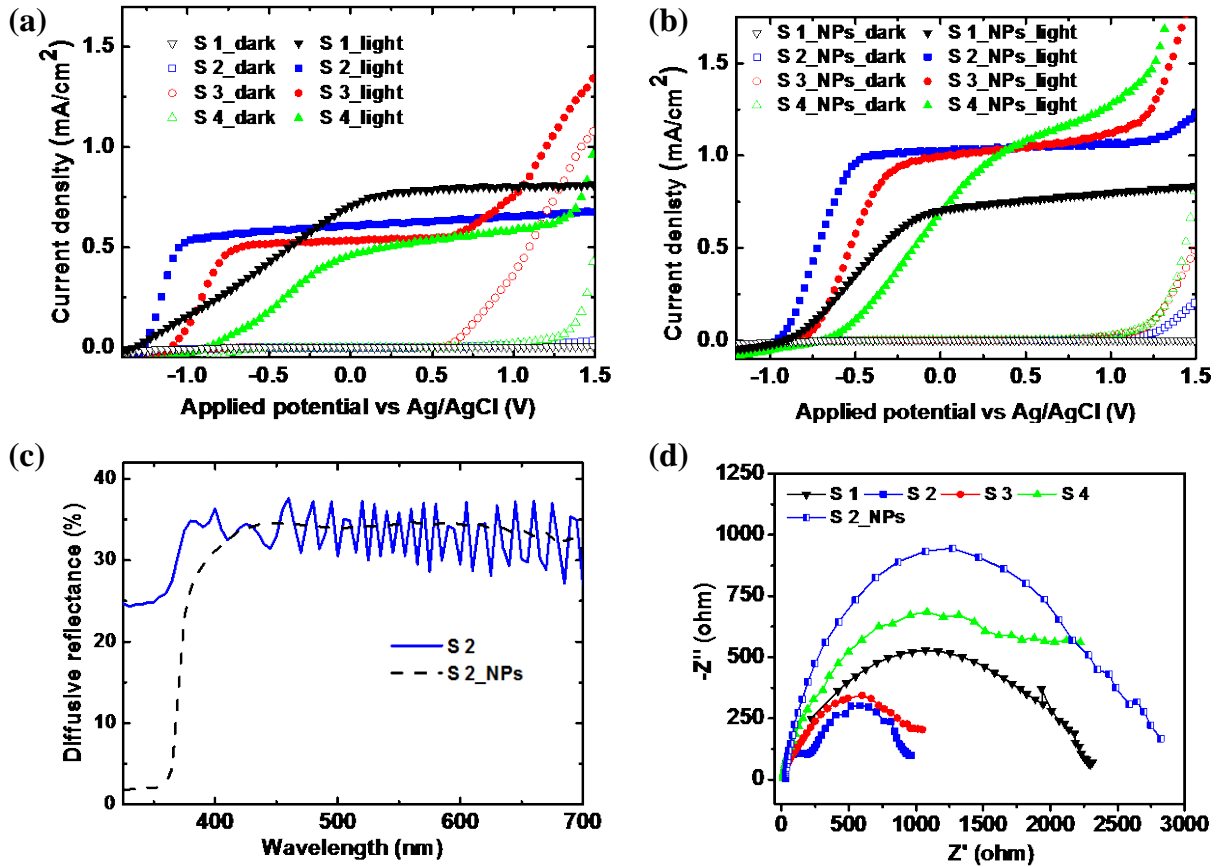


Fig 3. Linear sweep voltammetry of planar GaN photoanodes (a) and GaN NP photoanodes (b), diffuse reflectance spectra (c) and Nyquist plots (d).

Conductivity of the samples plays a critical role in the photoresponse of the photoanodes. As listed in table 1, GaN epilayers were grown with four different carrier concentrations and were used to fabricate NPs using top-down approach. PEC performance of the planar GaN and their respective NPs were evaluated by applying them as a photoanode in a three electrode configuration using a potentiostat. Figure 3(a) shows the PEC performance of GaN planar photoanodes with

varying carrier concentration. The dark currents for all the samples used in this study, irrespective of the carrier concentration, are insignificantly low. On the other hand, photoresponse of the photoanodes shows a strong dependence on carrier concentration. The onset potential is the lowest for S1 and increases with increasing carrier concentration. However, the photocurrent of S1 increases gradually and reaches a plateau at a much higher potential, indicating a large overpotential is needed for the undoped sample. The lowest overpotential to reach saturation currents is achieved for the photoanode S2 and it increases with increasing carrier concentration. This is attributed to the variation of depletion layer thickness with carrier concentration.[39] In general, the energy bands of the semiconductor are flat in bulk and bend at the surface. The band bending at the surface leads to the formation of charge neutral region (depletion region) near the semiconductor surface. Photogenerated carriers within the depletion zone can be effectively separated due to band bending as compared to those generated in the bulk region due to the flat nature of the bands in the bulk region of the semiconductor. To avoid the inevitable carrier recombination and extract the photogenerated carriers outside the depletion region, a higher overpotential is required. Typically, the thickness of the depletion layer is close to 100 nm for undoped samples[40] and decreases with increasing carrier concentration (Table S11). It is noticeable that the overpotential increases with increasing carrier concentration for S2, S3 and S4. However, S1 exhibits the highest overpotential despite the lowest carrier concentration. This is attributed to the poor conductivity of the undoped sample. On the other hand, onset potentials increases with increasing carrier concentration due to decreasing depletion width and the incorporation of structural defects in the epilayer with increasing doping content which may act as recombination centres for the photogenerated carriers.[41] It can be concluded that the lowest doped S2 photoanode has an optimum carrier concentration for lower carrier recombination and higher carrier separation efficiency.

PEC performance of GaN NP photoanodes with different carrier concentrations, fabricated from 120 s ICP-RIE etching of 5 nm Ni masked samples, is shown in Fig 3(b). The dark current for NPs photoanodes remains very low. Substantial improvement in saturation photocurrent is noticed for all the doped NP photoanodes compared to the planar ones. The highest saturation current density of around 1.0 mA/cm² is measured for all the NP photoanodes. To our knowledge, this is the best reported current density in the literature for GaN NP photoanodes without InGaN quantum wells under one sun light illumination which is also close to the theoretical limit.[42, 43]

This enhancement can be attributed to the enhanced absorption by the NP photoanodes apart from the increased semiconductor/electrolyte interface area. Fig 3(c) shows the diffuse reflectance spectra of the NP photoanode and its planar counterpart. It clearly shows that NPs absorb most of the light that is incident compared to the planar structures which reflect nearly 24% of the incident light near the band edge of GaN. This enhancement in absorption contributes to the generation of additional carriers in the NPs that participate in solar water splitting. Further, NPs offer massive semiconductor/interface area compared to the planar structures that result in enhanced depletion region around the semiconductor surface. This increased depletion region leads to the enhancement in carrier separation and efficiently drives the holes to the semiconductor-electrolyte interface.

Further, the onset and saturation potential shifts slightly towards higher potential for all the NP photoanodes irrespective of the carrier concentration with respect to the planar photoanodes. This anodic shift is attributed to the surface damage caused to the NP surface during plasma etching.[38] The surface defects act as traps for photogenerated charge carriers at the NP surface and therefore higher potential is required to extract the photogenerated carriers that are trapped in these defects. This leads to higher onset potential for NP photoanodes compared to the planar photoanodes. Electrochemical impedance measurements were carried out to understand the charge transfer kinetics between semiconductor and electrolyte as shown in Fig 3(d). The diameter of the half circles in the Nyquist plot is linked to the charge transfer resistance between semiconductor surface and electrolyte. Typically, the smaller the diameter the lower is the charge transfer resistance. The Nyquist plots indicate that the charge transfer resistance is increased for NPs photoanodes compared to the planar photoanodes. The increased charge transfer resistance is due to plasma etching damage for NP photoanodes and contributes to anodic shift in the onset potential and over potentials. Among the planar photoanodes, S2 exhibits the lowest charge transfer resistance. For photoanodes with higher carrier concentration, structural defects created by high doping and reduced depletion region causes increased charge transfer resistance between the semiconductor surface and electrolyte.

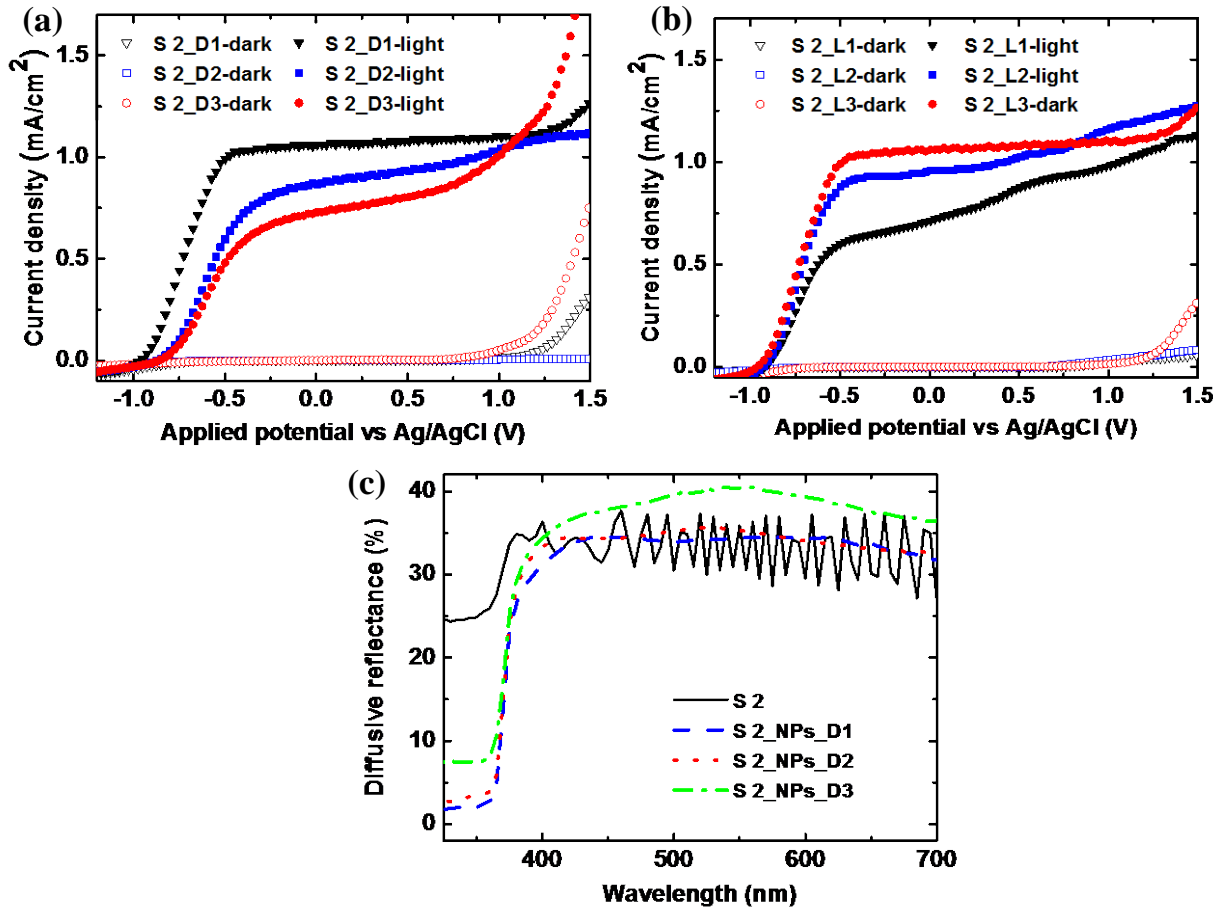


Fig 4. Linear sweep voltammetry of GaN NP photoanodes of varying diameters (a), varying lengths (b) and diffuse reflectance of NP photoanodes with varying diameters (c).

To investigate the effect of diameter and length on the PEC performance of NP photoanodes, NPs with different diameters and lengths are fabricated using S2. The diameters of the NP photoanodes are classified as a small (D1), medium (D2) and large (D3) fabricated from 5, 10 and 15 nm Ni masked samples, respectively. The diameter of the NPs is increased while the density is decreased with increasing Ni mask thickness under the same processing conditions as observed in Fig 2 (d) to (f). As shown in Fig 4(a), smallest diameter NP photoanodes exhibit the highest saturation photocurrent density and it decreases with increasing NP diameter. Increasing the diameter of NPs results in the reduction of semiconductor/electrolyte interface area. Consequently the depletion region is reduced with increasing NP diameter and causes inefficient separation of photogenerated carriers which could in turn contribute to the decreased saturation current density. In order to understand the effect of NP diameter on light absorption, diffusive

reflectance measurements were carried out and the results are shown in Fig 4(c). Reflection losses for NPs are significantly lower than epilayers irrespective of the diameters of the NPs. However, the reflection losses slightly increased with increasing diameter of the NPs. Higher reflection losses leads to the reduced light absorptions and carrier generation. Further, from Fig 4(a), it can also be observed that there is a slight change in onset potential with NP diameter. Due to the presence of surface band bending, the NPs could be either completely or partially depleted depending on the diameter of NPs for a given carrier concentration. The average diameter of the nanopillars is ~84, 144 and 202 nm for D1, D2 and D3 NP photoanodes, respectively. The estimated depletion layer width in one dimension is ~24.5 nm for $1 \times 10^{18} \text{ cm}^{-3}$ carrier concentration (Table SII).[44] For a given carrier concentration, the width of the depletion region is constant. Therefore, width of the flat band region increases with increasing NP diameter [45]. The deduced thickness of the flat band region is ~35, 95 and 153 nm for D1, D2 and D3 NP photoanodes, respectively. As a consequence, recombination of the photogenerated carriers becomes severe with increasing diameter of the NPs which could in turn contribute to increased onset potential with increasing NP diameter. However, the role of depletion region may become insignificant compared to the flat band region beyond certain NP diameter.[45]

Figure 4 (b) shows the influence of NP length on the current density of NP photoanodes under dark and light illumination. NPs fabricated from 5 nm Ni mask pattern were chosen for this study and the lengths of the NPs were 200, 500 and 900 nm labelled as short (L1), medium (L2) and long (L3) NP photoanodes, respectively. Since the NPs diameter is fixed while studying the influence of NPs length on PEC performance, insignificant changes in onset potential is observed regardless of the NP length. The onset potential is dependent on the depletion layer width that varies only with the diameter and doping concentration. On the other hand the saturation current is improved with increasing NP length. NPs/electrolyte interface area is estimated to be linearly dependent on the length of the NPs which contributes to increase depletion region volume with increasing NPs length.[44] The increase in depletion region volume leads to increased saturation level of the photo current. It is also expected that a certain minimum length of NPs is necessary to mitigate the optical reflection losses.[46] This could be another possible reason for the reduction in photocurrent density with reducing length of the NPs.

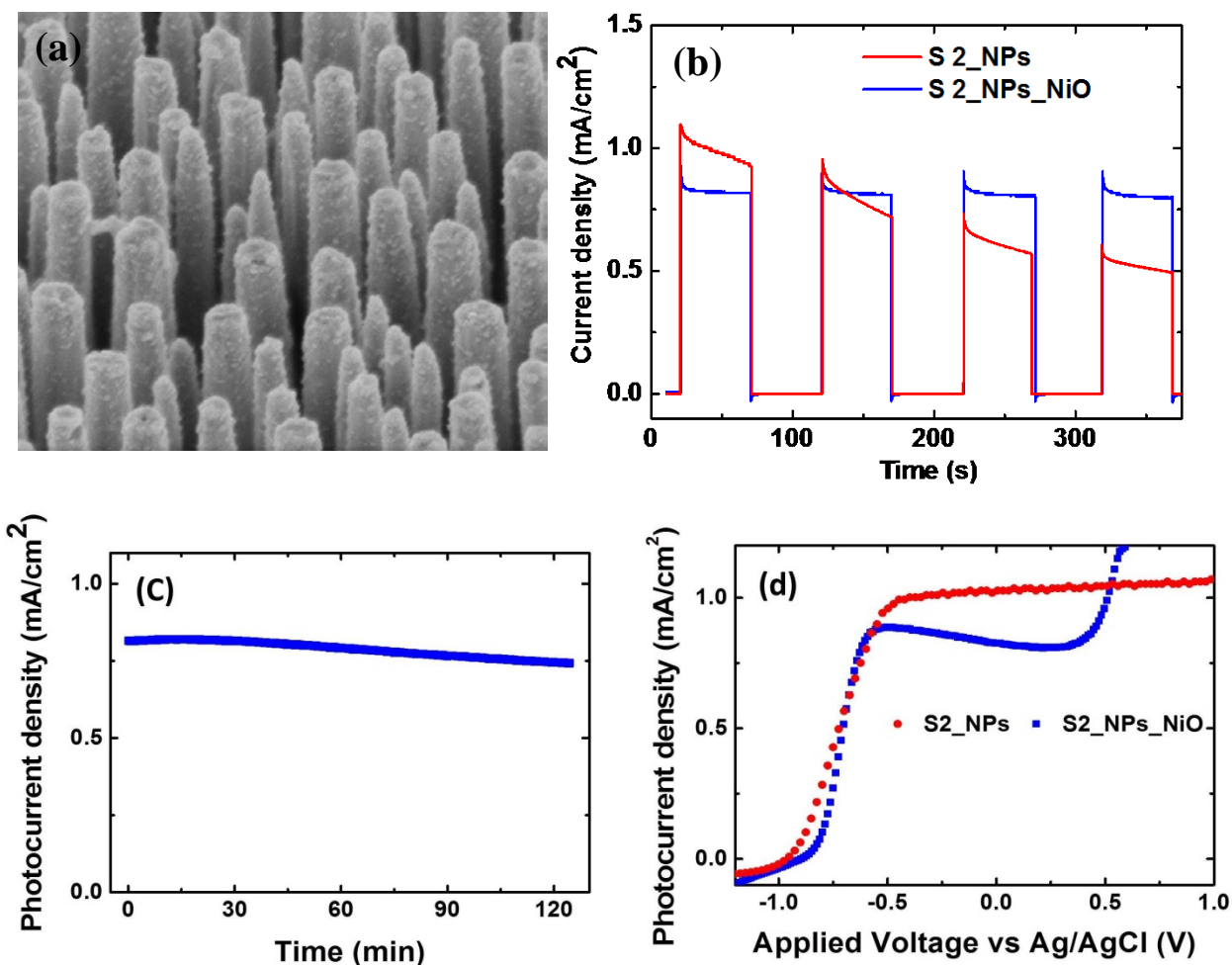


Fig 5. (a) SEM image of NiO co-catalyst deposited GaN NPs, (b) short term photostability of NP photoanodes with and without NiO co-catalyst and (c) long term photostability of NiO coated NP photoanodes under AM1.5 light illumination at 0 V vs. Ag/AgCl in 1M NaOH electrolyte and (d) linear sweep voltammetry of GaN NP photoanodes with and without NiO co-catalyst

It is well known that c-plane of GaN is chemically stable compared to the m-plane. [47, 48] When the NP photoanodes are dipped in electrolyte solution the m-plane is exposed to the electrolyte and results in photocurrent degradation with time. In order to improve the photostability of the NPs, NiO co-catalyst was coated around NP surface. From the SEM image shown in Fig 5(a), it can be observed that NiO co-catalyst nanoparticles are uniformly decorated around the NPs. The transient photocurrent curves for NP photoanodes with and without NiO co-catalyst are shown in Fig 5(b). The photocurrent density became relatively stable after NiO coating. The improved stability can be attributed to discontinuities in the bands between GaN and NiO [36, 41, 49]. NiO energy band positions are located at higher energy relative to the GaN energy bands (Fig.

SI2). As a consequence, the migration/extraction of photogenerated holes from GaN to NiO is enhanced, thereby significantly reducing the number of holes accumulated at the NPs/electrolyte interface. The reduced number of accumulated holes at the interface thus leads to reduced oxidative etching of the NPs surface. The long term stability of the NiO coated NP photoanodes is shown in Fig 5(c). Slight drop in the photocurrent density is noticed after 2hr of illumination. Figure 5(d) shows the linear sweep voltammetry of NP photoanodes with and without NiO catalyst under solar illumination. Small drop in the photocurrent density is observed for NiO coated NP photoanodes which could be due to the parasitic light absorption by the NiO nanoparticles. Further, a sharp rise in photocurrent density is observed for NiO coated GaN NP photoanodes beyond 0.5 V. This voltage is close to the reported OER potential of NiO catalyst[50] and hence the sharp raise of photocurrent density beyond this value can be attributed to the electrochemical catalytic activity of NiO.

Conclusions

To conclude, GaN NP photoanodes were fabricated using the top-down approach and exhibited substantial enhancement in the PEC performance compared to their planar counterparts. This enhancement is found to be strongly influenced by the carrier concentration and nanopillars dimensions. The photocurrent of n-doped NP photoanodes can reach up to $1\text{mA}/\text{cm}^2$, which is close to the theoretical limit, owing to their enhanced optical absorption and increased semiconductor/electrolyte interface. An anodic potential shift for the NP photoanodes is observed with increasing carrier concentration due to decreased depletion layer thickness and increased structural defects. Photocurrent density is reduced with increasing diameter and reducing length of the NPs. Reduction in optical absorption and depletion region can also contribute to the drop in the photocurrent density. Furthermore, increasing the diameter of the NP photoanodes for a given carrier concentration is found to cause an anodic shift of the onset potential due to the increased flat band region of the NPs. Finally, photodegradation of the NP photoanodes is substantially reduced by decorating the GaN NPs with NiO co-catalysts.

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Supporting Information

Improved Photoelectrochemical Performance of GaN Nanopillar Photoanodes

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Calculation of depletion layer width:

$$\text{Depletion layer width } (W) = \sqrt{\frac{\epsilon_0 \epsilon \Delta V}{e N_d}}$$

Where ϵ_0 = Vacuum permittivity

ϵ = dielectric constant of GaN

ΔV = potential difference between GaN and the redox couple in the solution

e = elementary charge

N_d = Carrier concentration

Table SI1: Estimated depletion layer width with different carrier concentrations in GaN epilayers

Carrier concentration (cm ⁻³)	Depletion layer width (nm)
1X10 ¹⁸	24.5
5X10 ¹⁸	11
1X10 ¹⁹	7.7

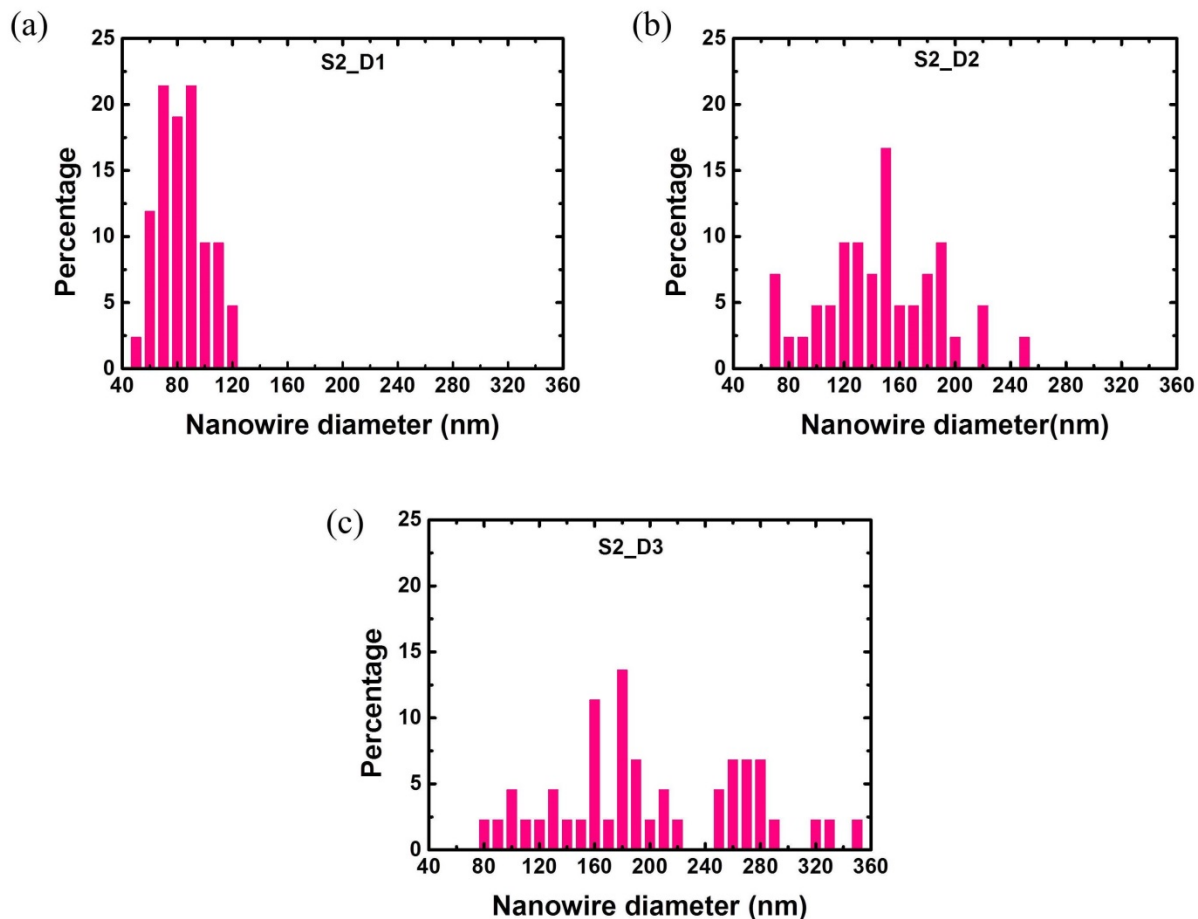


Figure SI1: Distribution of nanopillars diameters fabricated using (a) 5, (b) 10 and (c) 15nm Ni film.

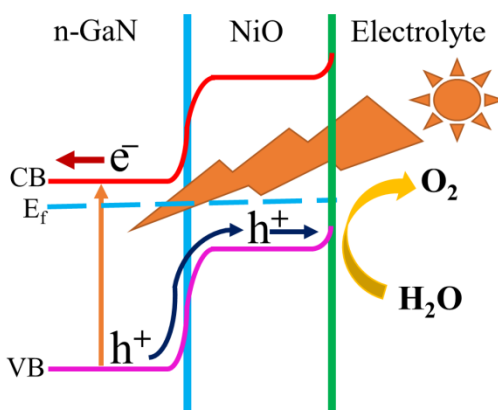


Figure SI2: Schematic of energy band diagram of GaN/NiO/electrolyte and electron and hole transfer at the material interfaces (CB: conduction band, VB: valency band and E_f : Fermi level).